

Reflectivity in Dense Xe Plasma

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Reflectivity can be used as a diagnostic tool for the investigation of dense plasma and its metallization. Measurements in shock wave produced dense Xenon plasmas at a temperature of around 30 000 K have been performed with a laser beam of three wavelengths: 1.06 μm , 0.694 μm and 0.532 μm [1]. A strong increase of the reflectivity indicating an onset of metallic behaviour was observed for densities above the critical density $n_e^{\text{cr}} = \epsilon_0 m_e \omega_L^2 / e^2$, where the plasma frequency coincides with the laser frequency.

Assuming a step-like shock front, the reflectivity is directly related to the dielectric function via the Fresnel formula [2]. The dielectric function is calculated from the generalized Drude formula

$$\epsilon(\omega) = 1 - \frac{\omega_{\text{pl}}^2}{\omega[\omega + i\nu(\omega)]}. \quad (1)$$

The collision frequency $\nu(\omega)$ was determined from a frequency dependent dynamical collision frequency for electron collisions taking into account dynamical screening and strong collisions. We also considered electron-electron collisions via a renormalization factor as well as electron-atom collisions. Bound state transitions contribute to the dielectric function following a cluster expansion. Despite these consistent approximations, the results for the collision frequency overestimate the reflectivity. Note, that molecular dynamics simulations also give an overestimation. In conclusion, a theoretical approach assuming a step like shock wave front is not able to explain the low reflectivity even above the critical density.

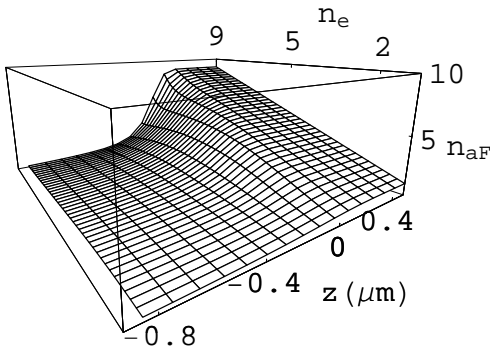


Figure 1: Adjusted density profile $n_{\text{aF}}(n_e, z)$ of the shock front, n_e and n_{aF} in 10^{21}cm^{-3}

This suggests more general assumptions about the spatial profile of the free electron density n_e within the shock wave front. In this case, the reflectivity is calculated using the solution of the Helmholtz equation

$$\frac{d^2 E(\omega_L, z)}{dz^2} + \frac{\omega_L^2}{c^2} \epsilon(\omega_L, z) E(\omega_L, z) = 0 \quad (2)$$

for the electric field. Using the measurements for the laser wavelengths 1.06 μm and 0.694 μm , we constructed a spatial

profile $n_{\text{aF}}(n_e, z)$ for every final value of n_e in the produced plasma with z the distance from the high density region. An asymmetric Fermi profile [3]

$$n_{\text{aF}}(n_e, z) = \frac{n_e}{e^{Y(z)} + 1}, \quad (3)$$

$$Y(z) = -\frac{z}{A} - e^{z/B} + C. \quad (4)$$

is deduced and illustrated in Fig. 1. We find a profile width of about 1 μm . Using this profile, we calculated the reflectivity for the third laser wavelength. Fig. 2 illustrates the experimental results in comparison to the theoretical model of an extended shock wave front. The good agreement supports our assumption of a density profile during the plasma creation.

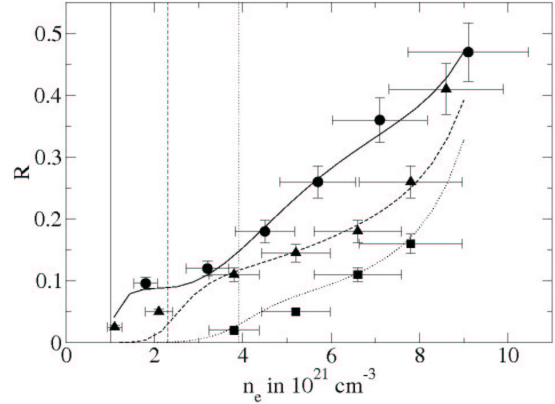


Figure 2: Reflectivity R as function of electron density; experimental results for 1.06 μm (\bullet), 0.694 μm (\blacktriangle) and 0.532 μm (\blacksquare); in comparison to calculation for asymmetric Fermi profile; vertical lines indicate the critical density of the laser wavelengths

A microscopic approach may be found on the basis of simulations. First attempts using hydrocodes fail to explain this empirical approach. A kinetic approach taking into account nonequilibrium situations might have to be adapted to the experimental situation. Work is in progress.

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References

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